

# Role of electron concentration parameter $e/a$ in energy-gap formation mechanism through interference phenomenon

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Dr.E.Belin-Ferré and I met for the first time at LAM8 conference in Vienna in 1992. We learned that both of us share a common interest in the electronic structure and electron concentration parameter  $e/a$  of complex compounds and could produce successive joint papers over 1994 to 2000. After the retirement from Nagoya University in 2005, I have been engaged in studying long-standing  $e/a$  issues for transition metal (TM) compounds in collaboration with my co-worker Prof.H.Sato, with the aim at deepening a theoretical insight into the Hume-Rothery electron concentration rule for TM compounds.

Mizutani and Sato have developed the *FLAPW-Fourier theory*, which is capable of making *ab initio* determination of the number of itinerant electrons per atom,  $e/a$ , for elements and compounds, almost regardless of the degree of metallicity, covalency and ionicity involved in the van Arkel-Ketelaar triangle diagram [1-3]. This certainly promises us to pave the way to establish the theoretical basis of the electron theory for any compounds involving TM and/or rare earth elements.

The theory makes a full use of the formalism of the FLAPW (Full-potential Linearized Augmented Plane Wave) electronic structure calculations based on the density functional theory: the  $j$ -th wave function in the interstitial region with the energy eigenvalue  $E^j$  at the wave vector  $\mathbf{k}_i$  obtained by partitioning the first Brillouin zone into  $N_{\mathbf{k}}$  meshes, is expanded into a Fourier series:

$$y^j(\mathbf{r}, \mathbf{k}_i) = \frac{1}{\sqrt{V}} \sum_p \hat{a}_j C_{\mathbf{k}_i + \mathbf{G}_p}^j \exp\{i(\mathbf{k}_i + \mathbf{G}_p) \cdot \mathbf{r}\}, \quad (1)$$

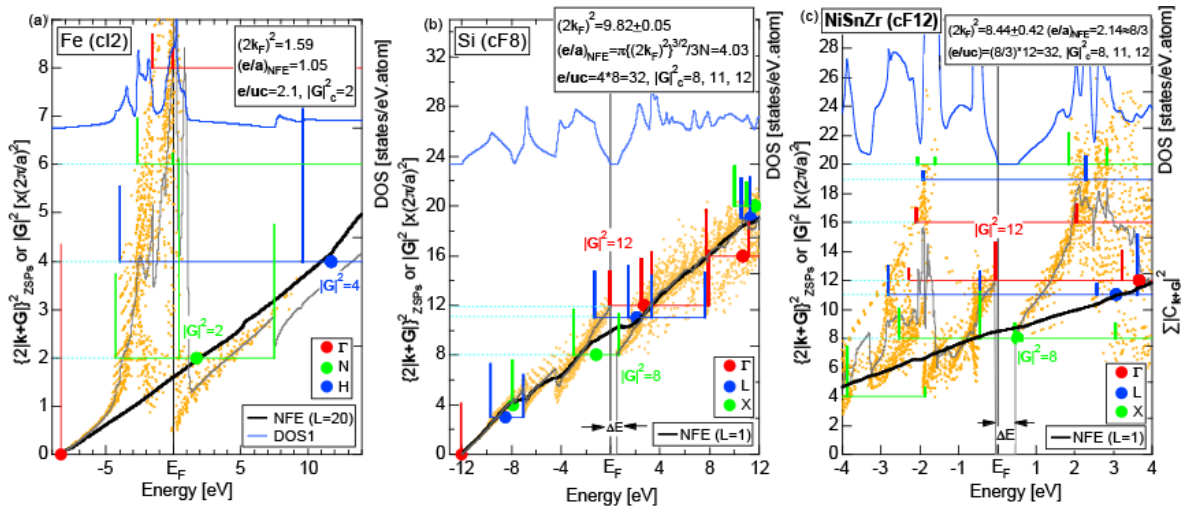
where  $V$  is the volume of the unit cell and  $\mathbf{G}_p$  is the reciprocal lattice vector of a given system. The square of the Fourier coefficient  $\left| C_{\mathbf{k}_i + \mathbf{G}_p}^j \right|^2$  forms a matrix with the electronic state  $\left| 2(\mathbf{k}_i + \mathbf{G}_p) \right|^2$  in row and energy eigenvalue  $E^j$  in column. The Fourier coefficients are plotted on the diagram with  $\left| 2(\mathbf{k}_i + \mathbf{G}_p) \right|^2$  as ordinates and  $E^j$  as abscissa with the choice of  $\mathbf{k}_i$  in two ways: one all  $\mathbf{k}_i$ 's in the Brillouin zone and the other  $\mathbf{k}_i$ 's only on its high-symmetry points. The center of gravity energy  $E_{\mathbf{k}_i + \mathbf{G}_p}^{cg}$  is calculated for each electronic state  $\left| 2(\mathbf{k}_i + \mathbf{G}_p) \right|^2$  by using the relation:

$$E_{\mathbf{k}_i + \mathbf{G}_p}^{cg} = \frac{\sum_j E^j(\mathbf{k}_i) \left| C_{\mathbf{k}_i + \mathbf{G}_p}^j \right|^2}{\sum_j \left| C_{\mathbf{k}_i + \mathbf{G}_p}^j \right|^2}, \quad (2)$$

where the sum is taken in the state  $|\mathbf{k}_i + \mathbf{G}_p\rangle$  over first  $L$  Fourier coefficients in the descending order in the  $j$ -th wave function.

The set of  $\frac{\hbar}{c} \left| 2(\mathbf{k}_i + \mathbf{G}_p) \right|^2, E_{\mathbf{k}_i + \mathbf{G}_p}^{cg}$  represents the *NFE* dispersion relation of itinerant electrons for any crystals including TM elements and their compounds. The  $\mathbf{e}/\mathbf{a}$  and  $\mathbf{e}/\mathbf{uc}$ , the product of  $\mathbf{e}/\mathbf{a}$  and the number of atoms per unit cell, can be calculated from  $(2k_F)^2$  read off from the intercept of the *NFE* line with the Fermi level. The *critical* reciprocal lattice vector  $|\mathbf{G}_c^2| \circ \left| 2(\mathbf{k}_i + \mathbf{G}_p) \right|_{ZSPs}^2$  responsible for yielding a gap through the interference phenomenon can be extracted from  $E_{\mathbf{k}_i + \mathbf{G}_p}^{cg}$  closest to the Fermi level.

Figure 1 shows the *FLAPW-Fourier spectra* for  $\alpha$ -Fe (cI2) in non-magnetic states, Si (cF8) and the half-Heusler compound NiSnZr (cF12): yellow dots represent the distribution of finite Fourier coefficients at  $E^j$  with all possible combinations of  $\mathbf{k}_i$  and  $\mathbf{G}_p$ . The coloured vertical line segments represent the maximum Fourier coefficients on high-symmetry points and are drawn in proportion to their magnitudes. The coloured circles indicate  $E_{\mathbf{k}_i + \mathbf{G}_p}^{cg}$  on the high-symmetry points. From the intersection of the *NFE* line with the Fermi level, we obtain  $(2k_F)^2 = 1.59 \pm 0.16$ ,  $\mathbf{e}/\mathbf{a} = 1.05$ ,  $\mathbf{e}/\mathbf{uc} = 2.1$  and  $|\mathbf{G}_c^2| = 2$  for  $\alpha$ -Fe. The data for both Si and NiSnZr are incorporated in Figure 1. The possession of common  $\mathbf{e}/\mathbf{uc} = 32$  and  $|\mathbf{G}_c^2| = 8, 11$  and 12 for both Si and NiSnZr is consistent with the fulfilment of the interference condition  $(2k_F)^2 = |\mathbf{G}_c^2|$  or  $\mathbf{e}/\mathbf{uc} = \frac{\rho}{3} \left\{ |\mathbf{G}_c^2| \right\}^{3/2}$  and explains the origin of the energy gap in terms of the *common* interference phenomena of itinerant electrons with set of zone planes associated with  $|\mathbf{G}_c^2| = 8, 11$  and 12 [3].



**Figure 1.** FLAPW-Fourier spectra of (a)  $\alpha$ -Fe (cI2), (b) Si (cF8) and (c) NiSnZr (cF12) [3].

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3. U.Mizutani and H.Sato, *Philos.Mag.*, (2018), **98**, 1307.

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